4th Northern germany OpenFoam User meetiNg 2016

3D Simulation of an Argon magnetoplasmadynamic thruster with coaxial induced magnetic field

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Braunschweig, 29 September 2016

Introduction	Phys. model	The density-based method	Validation	Application to MPDT	Conclusion	App en dix

- 2 Physical model and Governing equations
- Proposed density-based numerical method for MHD flow
- 4 Verification of the proposed schemes
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What is Magnetoplasmadynamic thruster?

Definition

A Magnetoplasmadynamic thruster (MPDT) is a form of electrically powered spacecraft propulsion which uses the Lorentz force to generate thrust.







Schematic view of a self-field MPDT



MPDT in operation²

¹National Aeronautics and Space Administration

²NASA Facts, Glenn Research Center

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Motivation



Physical problems

- Real gas effects
- Thermal nonequilibrium
- Problem of friction
- Difficulty to seperate flow and discharge

Numerical difficulties

- The necessary coupling of partial differential equations systems (Elliptical and hyperbolic)
- Nonlinearities



- Calculate a plasma flow in the self-field MPDT by using a Central-Upwind scheme
- Obtain insight into the physics of thrust production and Energy dissipation

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Physics	al mode					

The model ...

we present is dedicated to the description of an electrically conducting but electrically neutral fluid.



We assumed that:

- The propellant gas (Argon) is injected into the discharge chamber as fully-ionized fluid
- The plasma flow is in a state of thermal equilibrium $T \approx T_e \approx T_i$
- Electrical sheat, Hall effect and radiation processes are neglected.



Governing equations

For compressible MHD flow:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \rho \mathbf{U} = \mathbf{0}$$
$$\frac{\partial \rho \mathbf{U}}{\partial t} + \nabla \cdot \left[\rho \mathbf{U} \mathbf{U} + \left(p + \frac{B^2}{2\mu_0} \right) \mathbf{I} - \frac{\mathbf{B}\mathbf{B}}{\mu_0} = \nabla \cdot \tau_{visc}$$

$$\frac{\partial \mathbf{B}}{\partial t} + \nabla \cdot (\mathbf{U}\mathbf{B} - \mathbf{B}\mathbf{U}) = -\frac{1}{\mu_0 \sigma} \Delta \mathbf{B}$$
$$\frac{\partial \rho E}{\partial t} + \nabla \cdot \left[\left(\rho E + \rho + \frac{B^2}{2\mu_0} \right) \mathbf{U} - \frac{\mathbf{B}\mathbf{B}}{\mu_0} \right] = \nabla \cdot \left[k_{th} \nabla T - \left(\frac{\eta \mathbf{J} \times \mathbf{B}}{\mu_0} \right) \right]$$

 ρ , **U** and **B**: are the average density of all species, the velocity and magnetic field vectors $\rho E = \frac{p}{\gamma - 1} + \frac{1}{2}\rho U^2 + \frac{B^2}{2m}$ is the total energy density of the plasma _

Permeability of the free space μ_0 , Curent density J, Electrical conductivity σ , Thermal conductivity coefficient k_{th}

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Governing equations

 $\sigma = 1.53 \times 10$

Electrical conductivity ...

is deducted from the Spitzer-Harm formulation ^a:

^aSpitzer, L. and Härm, R.: Transport phenomena in a completely ionized gas. Phys. Rev., 89, 977 (1953)

$$-2\frac{T^{\frac{3}{2}}}{\ln\Lambda} \qquad \ln\Lambda = \ln[\frac{12\sqrt{2}\pi(k_{B}\epsilon_{0}T)^{\frac{3}{2}}}{a^{3}n^{\frac{1}{2}}}]$$

The divergence Cleaning method ...

is coupled with the induction equation to ensure that $div(B) = 0^{-a}$

^aDedner, A., Kemm, F., Munz, C.D., Schnitzer, T., and Wesenberg, M.:*Hyperbolic Divergence Cleaning for the MHD Equations.* J. Comput. Phys., 175, 645-673 (2002)

$$\frac{\partial \mathbf{B}}{\partial t} + \nabla \cdot (\mathbf{U}\mathbf{B} - \mathbf{B}\mathbf{U}) + \nabla \psi = -\frac{1}{\mu_0 \sigma} \Delta \mathbf{B} \qquad c_h = \frac{CFL}{\Delta t \times max(\frac{1}{h})}$$
$$\frac{\partial \psi}{\partial t} + c_h^2 \nabla \cdot \mathbf{B} = -\frac{c_h^2}{c_d^2} \psi \qquad c_d = \sqrt{-\Delta t \frac{c_h^2}{I_h(C_r)}}$$

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Summerized convective and diffusive fluxes $^{\rm 4}$ with divergence cleaning terms:

$$\sum_{f} \phi_{f} W_{f} = \sum_{f} [\alpha \phi_{f+} (W_{f+} + \kappa_{f+}) + (1 - \alpha) \phi_{f-} (W_{f-} + \kappa_{f-}) + \omega_{f} (W_{f-} - W_{f+}) + \frac{1}{2} (\Phi_{f-} + \Phi_{f+})]$$

$$W_{f} = \begin{pmatrix} \rho \\ \rho U_{x} \\ \rho U_{y} \\ \rho U_{y} \\ \rho U_{z} \\ B_{x} \\ B_{y} \\ B_{z} \\ \rho E + P_{over} \end{pmatrix}, \kappa_{f} = \begin{pmatrix} 0 \\ S_{x} P_{over} \\ S_{y} P_{over} \\ -b_{f} U_{x} \\ -b_{f} U_{z} \\ -b_{f} U_{z} \\ -b_{f} U_{z} \\ -b_{f} (U \cdot B) \end{pmatrix}, \omega_{f} = \begin{pmatrix} 0 \\ \omega_{f} \\ \omega_{f} \\ 0 \\ 0 \\ \omega_{f} \\ 0 \end{pmatrix}, \Phi_{f} = \begin{pmatrix} 0 \\ -B_{y} b_{f} \\ -B_{z} b_{f} \\ S_{y} \psi \\ S_{z} \psi \\ 0 \\ 0 \\ c_{h}^{2} b_{f} \end{pmatrix}$$

Central-upwind interpolation schema

⁴Kurganov, A., Noelle, S., Petrova, G.: Semi-discrete central-upwind schemes for hyperbolic conservation laws and Hamilton-Jacobi equations. SIAM J. Sci. Comput., 23, 707-740 (2001).



Convective Terms
$$\alpha = \begin{cases} \frac{1}{2} & \text{for the KI method} \\ \frac{\Psi_{f+}}{\Psi_{f+} + \Psi_{f-}} & \text{for the KNP method} \end{cases}$$
Diffusive Terms
$$\omega_f = \begin{cases} \alpha \max(\Psi_{f+}, \Psi_{f-}) & \text{for the KT method} \\ \alpha(1-\alpha)(\Psi_{f+} + \Psi_{f-}) & \text{for the KNP method} \end{cases}$$

$$\Psi_{f+} = \max(c_f | S_f | + \phi_{f+}, c_f | S_f | + \phi_{f-}, 0)$$

(1

Central Method fluxes ⁵

$$\Psi_{f-} = min(c_f|S_f| - \phi_{f+}, c_f|S_f| - \phi_{f-}, 0)$$

 $c_f = min(c_+, c_-)$

d

Effective speed of sound $c_{\pm} = (\frac{1}{2} [a_{\pm}^2 + \frac{B_{\pm}^2}{\mu_0 \rho_{\pm}} + \sqrt{(a_{\pm}^2 + \frac{B_{\pm}^2}{\mu_0 \rho_{\pm}})^2 - 4a_{\pm}^2 \frac{B_{n,\pm}^2}{\mu_0 \rho_{\pm}}}])^{\frac{1}{2}}$

⁵Greenshields, C.J., Weller, H.G., Gasparini, L., and Reese, J.M.: *Implementation of semi-discrete,* non-staggered central schemes in a collocated, polyhedral, finite volume framework, for high-speed viscous flows. Int. J. Numer. Meth. Fluids, 63, 1-21 (2010).

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The Shock-cloud interaction problem

$$(\rho, u_{\mathbf{x}}, u_{\mathbf{y}}, u_{\mathbf{z}}, p, B_{\mathbf{x}}, B_{\mathbf{y}}, B_{\mathbf{z}}) = \begin{cases} (3.86, 0, 0, 0, 167.34, 0, 2.18, -2.18) & \text{if } x < 0.6 \\ (1, -11.25, 0, 0, 1, 0, 0.56, -0.56) & \text{if } x \ge 0.6 \end{cases}$$



Initial conditions for density and the geometry used for the cloud-shock interaction test case ⁶ and density distribution on the $N = 800 \times 800$ grid at t = 0.06 s. Density in (kg/m^3)

⁶Xisto, C.M., Pascoa, J.C., Oliviera, P.J.: *A pressure-based high resolution numerical method for resistive MHD.* J. Comput. phys., 275, 323-345 (2014).

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Density



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Villani-H thruster: Case set-up



Input of the code:

Discharge Current I (A)

Temperature T(K)The propellant mass flow rate $\dot{m}(kg/s)$

⁷K. Sankaran, 2005



Mesh grid of the Villani-H thruster with about 1.25 millions cells

 $B_0 = \frac{\mu_0 I}{2\pi r}$ GroovyBC

$$p =
ho RT$$

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Boundary conditions

Variables	Inlet	Electrodes	Insulated walls	O ut let
U (m/s)	$\dot{\mathbf{m}} = 6.0(g/s)$	slip/non slip	non-slip	ZG
T (eV)	fixedValue 1.0	ZG	ZG	ZG
p (Pa)	ZG	ZG	ZG	WT
B (T)	TDBC	Conducting walls	fixedValue (0,0,0)	fixedValue $(0, 0, 0)$

Boundary conditions of the MPDT simulations

- ZG: Zero gradient
- WT: Wave transmissive
- TDBC: Time depending boundary condition

$$B = \begin{cases} 0, & \text{if } t < t_1. \\ B_0 \frac{t - t_1}{t_2 - t_1}, & \text{if } t_1 < t < t_2. \\ B_0, & \text{if } t > t_2 \end{cases}$$

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Results



(c) Initial conditions for density and the geometry used for the cloud-shock interaction test case



(d) Density distribution

Velocity, Mach number distribution (left), Temperature and pressure (right) with $\frac{l^2}{m} = 05 \times 10^9 \cdot \frac{A^2 s}{kg}$, $l_c = 0.264$ m, $r_c = 0.0095$ m, $r_a = 0.051$ m and $\frac{l_a}{r_a} = 04$ for MPDT02. Units: Velocity (m/s), Temperature (K), Pressure (pa) and Density (kg/m^3)

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Geometric scaling analysis

MPDT geometries and $\frac{l^2}{m}$ values for the numerical parameter study on the HLRN. r_a , l_a and l_c denote anode radius, anode length and cathode length respectively.

Case geometry	$\frac{l_a}{r_a}$	$\frac{I^2}{\dot{m}} \left[10^9 \cdot \frac{A^2 s}{kg} \right]$	l _c [m]
MPDT01	1	5.7	0.132
$r_c = 0.0095[m]$	2	11.2	0.264
$r_a = 0.025[m]$	3	18.4	
	4	25.6	
	5	48.4	
		60.2	
MPDT02	1	5.7	0.132
$r_c = 0.0095[m]$	2	11.2	0.264
$r_a = 0.051[m]$	3	18.4	
	4	25.6	
	5	48.4	
		60.2	
MPDT03	1	5.7	0.132
$r_c = 0.0181[m]$	2	11.2	0.264
$r_a = 0.051[m]$	3	18.4	
	4	25.6	
	5	48.4	
		60.2	





(e) Thrust

(f) Specific impulse and current voltage

Thrust, efficiency, plasma current voltage for MPDT01 with $\frac{l_a}{r_a} = 04$.

$\frac{l^2}{\dot{m}} \left[10^9 \cdot \frac{A^2 s}{kg} \right]$	$F_{EM}[N]$	Maecker[N]	F _{thermal} [N]	F _{total} [N]	$V_{plasma}[V]$	lsp[s]
60.2	32.86	36.06	24.54	57.4	39.238	975.13





Thrust function of aspect ratio for MPDT01 for $\frac{l^2}{\dot{m}} = 05.7 \times 10^9 \cdot \frac{A^2 s}{kg}$ and $\frac{l^2}{\dot{m}} = 11.2 \times 10^9 \cdot \frac{A^2 s}{kg}$.





(g) Thrust

(h) Specific impulse and current voltage

Thrust, efficiency, plasma current voltage for MPDT03 with $\frac{l_a}{r_a} = 04$.

$\frac{l^2}{\dot{m}} \left[10^9 \cdot \frac{A^2 s}{kg} \right]$	$F_{EM}[N]$	Maecker[N]	F _{thermal} [N]	F _{total} [N]	$V_{plasma}[V]$	lsp[s]
60.2	34.04	36.24	28.20	56.24	26.425	955







Thrust efficiency (%) for both MPDT01, MPDT02 and MPDT03 with $\frac{l_{a}}{r_{a}}=$ 04.

Case geometry	η_{effmax} [%]	η_{effmin} [%]
MPDT01	36.25	20.27
MPDT02	27.14	16.41
MPDT03	53.35	21.9

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Conclu	sion					

- The central-upwind schemes have been succefully extended to MHD equations.
- The solver has demonstrated capability to compute resistive plasma flows in simple geometries

What remain to be done?

• improve the physical model (by adding Hall effect, real gas effect, considering multiple fluid plasma flow, ...) in order to achieve a predominantly electromagnetic acceleration mode for all thruster configuration in more realistic scenario. Thank You for your Attention!

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 $j \times B$: Lorentz force of a volume element

$$\mathbf{j} \times \mathbf{B} = \frac{1}{\mu} (\nabla \times \mathbf{B}) \times \mathbf{B} = \underbrace{\frac{1}{\mu} (\mathbf{B} \cdot \nabla) \mathbf{B}}_{mag. Diffusion} - \underbrace{\nabla \left(\frac{B^2}{2\mu}\right)}_{mag. Pressure}$$

Maecker (1955) introduces an analytical expression of the electromagnetic thrust for coaxial self-field MPD thrusters, based on continuum plasma and ideal MPD approximation.

$$F_{Maecker} = rac{\mu_0}{4\pi} I^2 (ln rac{r_a}{r_c} + A)$$

where A is a dimensionless constant between 0 and 1. In this study, we considered A = 0.



The MHD shock tube problem: Magnetic shock wave

Initial conditions of the problem ⁸

- Density jump
- Pressure jump
- Cross-sheared magnetic unsteadiness

⁸Sod, G.A.: A survey of several finite difference methods for systems of nonlinear hyperbolic conservation laws. J. Comput. phys., 27, 1-31 (1978)





Comparaison of exact y-component of magnetic field (a)and density (b) profiles with numerical simulation results at $t = 0.1 \ s$





Comparaison of exact y-component of velocity (a) and pressure (b) profiles with numerical simulation results at t=0.1~s





(m) Thrust

(n) Specific impulse and current voltage

Thrust, efficiency, plasma current voltage for MPDT02 with $\frac{l_a}{r_a} = 04$.

$\frac{l^2}{\dot{m}} \left[10^9 \cdot \frac{A^2 s}{kg} \right]$	$F_{EM}[N]$	Maecker[N]	F _{thermal} [N]	F _{total} [N]	$V_{plasma}[V]$	lsp[s]
25.2	24.43	27.2	3.57	28	31.34	475